## Physics 5455 – Problem set 6 – Revised Mar 18

1. Reflection-free potentials. A quantum particle moves in the potential

$$V(x) = -\frac{\mu(\mu+1)\hbar^2}{2ma^2[\cosh(x/a)]^2}$$

for  $\mu > 0$ .

- (a) Find the scattering states  $(E \ge 0)$  as follows. First, set  $\psi_k(x) = e^{ikx}\phi_k(y)$  for  $E = \hbar^2 k^2/2m$  and y = x/a. Then, substitute  $z = \tanh y$ , and show that the resulting power series terminates if  $\mu$  is an integer.
- (b) For  $\mu = 1, 2$ , determine the scattering states  $(E \ge 0)$  explicitly.
- 2. Dirac delta function potential. Consider the time-independent Schrödinger equation with potential

$$V(x) = \lambda \delta(x)$$

(a) Assuming that the wavefunction itself is continuous at x = 0, show by integrating the Schrödinger equation that the first derivative has a discontinuity given by

$$\lim_{x \to 0^+} \psi'(x) - \lim_{x \to 0^-} \psi'(x) = \frac{2m}{\hbar^2} \lambda \psi(0)$$

- (b) Find the reflection and transmission coefficients for a plane wave scattering state (E > 0) incident from the left.
- (c) For  $\lambda < 0$ , find the normalized bound state wavefunction and its energy eigenvalue (E < 0).
- (d) Show that the pole in the (analytic continuation of the) transmission amplitude from part (b), coincides with the bound state energy computed in (c).
- 3. The Kronig-Penney model. In this problem we shall study a simple model of electrons in a one-dimensional crystal, interacting only weakly with the ion cores. Begin with Schrödinger's equation for a particle of mass m in one dimension:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} + V(x)\psi(x) = E\psi(x)$$

(a) First, show that

$$\frac{p^2}{2m}\varphi(p) + \int_{-\infty}^{\infty} \frac{dq}{2\pi\hbar} \tilde{V}(p-q)\varphi(q) = E\varphi(p)$$

where  $\tilde{V}$ ,  $\varphi$  are the Fourier transforms of V,  $\psi$ , respectively.

(b) Assume the potential is periodic with period a, described by delta functions:

$$V(x) = A \sum_{n=-\infty}^{\infty} \delta(x+na)$$

Show that the equation above reduces to

$$\left(\frac{p^2}{2m} - E\right)\varphi(p) + \frac{A}{a}\sum_{n=-\infty}^{\infty}\varphi\left(p + \frac{2\pi\hbar}{a}n\right) = 0$$

You will need to use the identity

$$\sum_{n=-\infty}^{\infty} e^{i(\omega-\omega')n} = 2\pi \sum_{n=-\infty}^{\infty} \delta(\omega-\omega'+2\pi n)$$

(c) Define

$$f(p) = \sum_{n=-\infty}^{\infty} \varphi\left(p + \frac{2\pi\hbar}{a}n\right)$$

Show that

$$\varphi\left(p + \frac{2\pi\hbar}{a}n\right) = -\frac{2m(A/a)f(p)}{(p + (2\pi\hbar/a)n)^2 - 2mE}$$

(d) Sum over all n to derive the constraint

$$\frac{a}{2mA} = -\sum_{n=-\infty}^{\infty} \left[ \left( p + \frac{2\pi\hbar}{a} n \right)^2 - 2mE \right]^{-1}$$

(e) Use the identity

$$\cot x = \sum_{n=-\infty}^{\infty} \frac{1}{x + n\pi}$$

and trigonometric identities to show that the constraint above can be rewritten as

$$\cos(pa/\hbar) = \cos(ka) + \left(\frac{mA}{k\hbar^2}\right)\sin(ka)$$

for

$$k^2 \; = \; \frac{2mE}{\hbar^2}$$

Note that the left-hand side can only take values in the range [-1, 1], whereas the right-hand side can take values outside that range. The allowed values of k come in segments known as *bands*, and the gaps between bands are known as *band gaps*.