Insights into Hydration Dynamics and Cooperative Interactions in Glycerol–Water Mixtures by Terahertz Dielectric Spectroscopy

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ABSTRACT: We report relaxation dynamics of glycerol–water mixtures as probed by megahertz-to-terahertz dielectric spectroscopy in a frequency range from 50 MHz to 0.5 THz at room temperature. The dielectric relaxation spectra reveal several polarization processes at the molecular level with different time constants and dielectric strengths, providing an understanding of the hydrogen-bonding network in glycerol–water mixtures. We have determined the structure of hydration shells around glycerol molecules and the dynamics of bound water as a function of glycerol concentration in solutions using the Debye relaxation model. The experimental results show the existence of a critical glycerol concentration of ~7.5 mol %, which is related to the number of water molecules in the hydration layer around a glycerol molecule. At higher glycerol concentrations, water molecules dispersed in a glycerol network become abundant and eventually dominate, and four distinct relaxation processes emerge in the mixtures. The relaxation dynamics and hydration structure in glycerol–water mixtures are further probed with molecular dynamics simulations, which confirm the physical picture revealed by the dielectric spectroscopy.

1. INTRODUCTION

Along with water, a variety of cosolvents play important roles in biological systems.1–5 The presence of cosolvents changes the behavior of water such as hydrogen-bonding network, dynamics, polar property, and spatial distribution.6 Cosolvents can stabilize the activity of an enzyme and the native structure of a protein,7 increase the aqueous solubility of a nonpolar drug by several orders of magnitude,8 and enhance the chemical stability of a substance.7 The study of the dynamics of these chemical biomolecules is indispensable to get a comprehensive perception of their conduct in aqueous solutions. Glycerol (C₃H₈O₃) is an important cosolvent in this context, which was a subject of numerous studies in molecular dynamics (MD) simulations8–13 and experiments.14–16 At ambient conditions, this trihydric alcohol with three hydroxyl groups is a colorless, sugar-like, highly viscous liquid. The high flexibility and viscosity of glycerol make it an important system in studies of the glass transition.17 Also, glycerol has been used to preserve proteins because of its cryoprotective properties18 and to stabilize enzyme activities.2,3 The ability of glycerol to form hydrogen bonds with water makes glycerol–water mixtures fascinating solutions for enhancing the solubility of several drugs.4,19 Thus, a comprehensive understanding of the hydration dynamics and cooperative interactions in glycerol–water mixtures is needed to help us understand the role of glycerol in these activities. The investigation of dynamics at the molecular scales in a complex liquid is a major challenge in physical chemistry and chemical physics. Glycerol and glycerol–water mixtures have been the subject of numerous investigations including MD simulations,10–12 broadband dielectric spectroscopy,14,17,20–22 nuclear magnetic resonance,23 infrared spectroscopy,10 and Raman spectroscopy.24 Although such a wide range of techniques has been employed to investigate the hydrogen-bonding dynamics of glycerol–water mixtures in different frequency ranges, the hydration dynamics in aqueous glycerol solutions is yet to be elucidated. The dielectric relaxation spectroscopy, which measures the rearrangement dynamics in a hydrogen-bond network, is a handy tool that can be used to advance our understanding of the hydration structure and dynamics in glycerol–water mixtures.

Recent developments in megahertz-to-terahertz spectroscopy provide us a possibility to conduct dielectric response measurements in a wide range of timescales to reveal hydration dynamics and cooperative interactions in glycerol–water mixtures. We have adopted this technique to investigate the structure and dynamics of hydration shells and the properties of water molecules interacting with proteins and micelles. The results enabled us to map out the physical behavior of different molecules in their aqueous solutions.25–27 Our spectrometer covers a large spectral range from megahertz-to-terahertz frequencies and has a significantly improved signal-to-noise ratio with high power, providing high accuracy measurements.28 In the present study, we focus on the nature of hydration dynamics and the molecular dynamics of glycerol in its aqueous solutions at the molecular level. From the complex dielectric response, we have explored the relaxation processes in these solutions that span a wide range of glycerol...
concentrations. The behavior of the hydrogen-bond network related to distinct relaxation times of bulk water, water molecules in a hydration layer around glycerol, and water confined in a glycerol network has been discussed. A critical value of glycerol concentration is identified beyond which water molecules confined in a glycerol network start to emerge. The existence of different types of water with distinct relaxation times is further confirmed with MD simulations. Finally, a unified physical picture is presented to help us understand the colligative properties of glycerol–water solutions.

2. EXPERIMENTAL METHODS

2.1. Materials. Glycerol (≥99.5%) with a molecular weight of 92.093 g/mol, purchased from Sigma-Aldrich (cat. no. 56-81-0), was used to prepare glycerol–water mixtures. The mixtures with the glycerol content from 5 to 50% volume percentage with an increment of 5% vol % were prepared from pure glycerol and deionized water (resistivity of 18.2 MΩ cm). Measurements on pure glycerol and water were also performed, and the results are included in our discussion. Table 1 shows the glycerol volume percentage (vol %), weight-by-weight ratio (w/w), and a conversion to the glycerol molar percentage ($x_{\text{glyc}}$) of our glycerol–water mixtures.

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<th>glycerol volume percentage (vol %)</th>
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<th>$\tau_2$ (ps)</th>
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Table 1. Glycerol–Water Mixtures Used for the MHz-to-THz Dielectric Spectroscopy Measurements and Dielectric Relaxation Times, $\tau_1$, $\tau_2$, and $\tau_3$ for These Mixtures at Various Concentrations at 25 °C in Which the Relaxation of Bulk Water, $\tau_4$, is 8.27 ± 0.35 ps, and the Dielectric Constant at Higher Frequencies, $\varepsilon_{\text{rel}}$, is 7.06 ± 0.72.
50 MHz to 0.5 THz at 25 °C. In the frequency range from 50 MHz to 0.5 GHz, we use an enhanced open-end probe (Agilent 85070E) and a vector network analyzer (Agilent PNA N5225A). The calibration of this system was performed under three standards including air, pure water, and mercury (short circuit). The complex dielectric response including the real (dielectric dispersion), \( \epsilon''_{\text{sol}}(\nu) \), and the imaginary (dielectric loss), \( \epsilon''_{\text{sol}}(\nu) \), components was evaluated using Agilent software with an accuracy of \( \Delta \varepsilon/\varepsilon = 0.05 \). In the frequency range from 60 GHz to 0.5 THz, the dielectric response of glycerol–water mixtures has been collected using a gigahertz-to-terahertz spectrometer based on the above vector network analyzer with frequency extenders from Virginia Diodes. Our spectrometer is capable of simultaneously measuring intensity and phase over a large effective dynamical range. The solutions were kept in a sample cell made of anodized aluminum at 25 °C, and the temperature was controlled with an accuracy of ±0.02 °C using a Lakeshore 336 temperature controller.

The complex dielectric response of a glycerol mixture can be expressed as a function of frequency, \( \nu \):

\[
\epsilon_{\text{sol}}^*(\nu) = \epsilon_{\text{sol}}'(\nu) + i\epsilon_{\text{sol}}''(\nu)
\]

Figure 1 shows the complex dielectric response spectra for the mixtures with different glycerol concentrations from 0 mol % (pure water) to ~20 mol %. The dielectric spectra of pure glycerol and pure water are also included as reference. When the concentration of glycerol is increased, the absorption of the mixture decreases dramatically and the maximum of the dielectric loss shifts significantly toward lower frequencies. These changes are expected to occur as a water molecule and a glycerol molecule have different molecular weights and polarities, which affect the orientational relaxation of the associated dipoles.

### 3. RESULTS AND DISCUSSION

The dielectric properties of an aqueous solution present a complex behavior, originating from different, and/or partially overlapping, polarization mechanisms. In order to determine the contribution of each polarization mechanism in a solution to its dielectric response, the data were analyzed by simultaneously fitting the measured real, \( \epsilon_{\text{sol}}'(\nu) \), and imaginary, \( \epsilon_{\text{sol}}''(\nu) \), components to the Debye relaxation model based on a sum of four individual contributions:

\[
\epsilon_{\text{sol}}^*(\nu) = \epsilon_{\infty} + \frac{\epsilon_1 - \epsilon_\infty}{1 + i2\pi\nu\tau_1} + \frac{\epsilon_2 - \epsilon_1}{1 + i2\pi\nu\tau_2} + \frac{\epsilon_3 - \epsilon_2}{1 + i2\pi\nu\tau_3} + \frac{\epsilon_\infty - \epsilon_3}{1 + i2\pi\nu\tau_4}
\]

where \( \Delta \varepsilon_1 = \epsilon_1 - \epsilon_\infty, \Delta \varepsilon_2 = \epsilon_2 - \epsilon_1, \Delta \varepsilon_3 = \epsilon_3 - \epsilon_2, \) and \( \Delta \varepsilon_4 = \epsilon_\infty - \epsilon_3 \) are the associated dielectric strengths, corresponding to four relaxation processes with different relaxation time constants. Deconvolution of the dielectric spectra including the real and imaginary components has been performed for the glycerol–water mixture with a molar percentage, \( x_{\text{glycerol}} \) of 19.69 mol %. The complex dielectric spectra in Figure 2 for this mixture indicate that the four relaxation processes are centered at 168 ± 18 MHz (\( \tau_1 \approx 945 \pm 115 \) ps), 1.8 ± 0.2 GHz (\( \tau_2 \approx 88 \pm 9 \) ps), 4.0 ± 0.7 GHz (\( \tau_3 \approx 39 \pm 8 \) ps), and 19.2 ± 0.8 GHz (\( \tau_4 \approx 8.27 \pm 0.35 \) ps). Corresponding to the four components of the imaginary part, the deconvolution for the real part has been shown in the inset of Figure 2. The dielectric constant at higher frequencies obtained from this fitting procedure, \( \epsilon_{\infty} = 7.06 \pm 0.72 \), is within experimental uncertainty of the value reported in the literature. The fitting also yields the values of the dielectric strengths, \( \Delta \varepsilon_1 = 1.35 \pm 0.05, \Delta \varepsilon_2 = 1.26 \pm 0.27, \Delta \varepsilon_3 = 44.47 \pm 2.10, \) and \( \Delta \varepsilon_4 = 8.95 \pm 0.45 \). The relaxation
The dielectric parameters of glycerol-water mixtures allow us to evaluate the relaxation times of molecules in the mixtures. Interestingly, the relaxation times for the four processes identified by fitting the dielectric relaxation spectra to the Debye model are almost constant when the glycerol concentration is varied. The slowest relaxation time, $\tau_1$, is around 910 ± 135 ps or 174 ± 22 MHz (Figure 3, inset). This mode is called the $\beta$-relaxation, which is a typical relaxation process in glass-forming liquids. This value is slightly faster than those in pure glycerol.17,35,36 Furthermore, the dielectric strength of this relaxation process increases approximately linearly with the glycerol concentration (Figure 3, inset). The indication is that glycerol molecules form strong hydrogen bonds with water molecules in an aqueous solution, and these bonds are similar to the hydrogen bonds between glycerol molecules in pure glycerol. As a result, glycerol in an aqueous solution exists in an environment which is not very different from that in pure glycerol as far as its rotational motion is concerned.

The dielectric parameters of the $\beta$-relaxation process, including the dielectric strength, $\Delta \varepsilon_{\beta}$, and the relaxation time, $\tau_1$, allow us to evaluate the hydrodynamic radius of a glycerol molecule and its electric dipole moment, $\mu$, in glycerol–water mixtures. In a simple physical model, a dipole is regarded as a sphere whose rotation in response to an electric field is opposed by the hydrodynamic friction with the surrounding solvent. The relaxation time of a spherical molecule in a dilute solution with hydrodynamic radius, $R$, rotating in a medium with macroscopic viscosity, $\theta$, is given by the Debye equation37

$$\tau_1 = \frac{4\pi R^3 \theta}{k_B T}$$

(3)

where $k_B$ is the Boltzmann constant, and $T$ is the temperature. At 25 °C, the viscosity of water is about 0.89 x 10^{-3} kg m^{-1} s^{-1} (Pa s).38 From the measured values of $\tau_1$, we derive the hydrodynamic radius of glycerol in low glycerol concentration solutions as $R = 7.06$ Å. Note that this radius is larger than the unit cell dimensions of a glycerol molecule ($a = 4.8 \text{ Å}$, $b = 5.1 \text{ Å}$, $c = 7.8 \text{ Å}$).38,39 This difference is expected as the hydrodynamic radius estimated from the dielectric measurements includes the contribution from the bound water molecules around a glycerol molecule.

The effective dipole moment of glycerol in an aqueous solution can be estimated from the dielectric strength. Several approximations for such estimation have been suggested. We adopted the same approach as in ref 40 to calculate the effective dipole moment of glycerol, $\mu_{\text{eff}}$, by using the Onsager–Onsley model33

$$\mu_{\text{eff}}^2 = \frac{2\varepsilon_0 k_B T \Delta \varepsilon_{\beta}}{N a g_{\text{ff}} c}$$

(4)

where $N_a$ is Avogadro’s number, $c$ is the molar concentration (mol/m^3) of glycerol in the solution, $\varepsilon_0$ is the permittivity of vacuum, and $g_{\text{ff}}$ denotes the Kirkwood correlation factor and is often assumed to be 1 in a diluted solution.37,42 The estimated values of the effective dipole moment of glycerol at various concentrations are shown in the lower inset of Figure 3. At a low glycerol concentration, $\mu_{\text{eff}}$ is roughly a constant of 0.85 D, and it starts to increase with an increasing glycerol concentration when the concentration is higher than about 7.5 mol %. At low concentration, each glycerol molecule is well covered by a hydration layer that is separated from other hydration layers. When the glycerol concentration increases beyond the threshold, glycerol molecules tend to form clusters or networks and hydration layers start to overlap. Because the dipole moment of glycerol is 2.56 D, an increase of the effective dipole moment of glycerol is observed as the glycerol concentration is increased further.

3.2. Bulk Water Relaxation and Hydration Effect. The dielectric spectroscopy provides insights into the dynamics of bulk water, water bound to glycerol, and water confined in a glycerol network. The relaxation time for bulk water, $\tau_2 \approx 8.27$ ps, is independent of glycerol concentration and similar to the values reported in the literature for pure water at gigahertz frequencies.32,33 However, when glycerol is added to a glycerol–water mixture, the dielectric strength, $\Delta \varepsilon_{\text{ff}}$, of the bulk water in the mixture (Figure 4a) decreases significantly. The lowering of the dielectric response with an increasing glycerol concentration comes from two main reasons. First, the presence of glycerol will reduce the concentration of water in the mixture, thus lowering the dielectric response of water. Second, water molecules form hydrogen bonds with glycerol. A glycerol molecule has three OH groups and can form six hydrogen bonds with surrounding water in its hydration layer. When the glycerol concentration is increased, the amount of water in the hydration layers also increases. These water molecules have a different relaxation process than the bulk
Debye model. Specifically, relaxation processes with a frequency of $4.5 \pm 0.9$ GHz (Figure 4b, inset) and another one at $1.87 \pm 0.22$ GHz (Figure 5, inset), corresponding to the time constants of $35 \pm 8$ and $85 \pm 9$ ps, respectively, are identified. These time constants are much longer than that of pure water, which is 8.27 ps. As confirmed later with MD simulations, the relaxation time, $\tau_2 \approx 35 \pm 8$ ps, can be associated with water molecules bound to a glycerol molecule and forming its hydration layer, whereas the even longer relaxation time, $\tau_3 \approx 85 \pm 9$ ps, can be assigned to water molecules confined in a glycerol network, that is, water molecules shared by more than one glycerol molecule. The longer relaxation times of the bound and confined water indicate that the hydrogen bond between a glycerol molecule and a water molecule is stronger than that between two water molecules.

The amount of bound water molecules that do not take part in the bulk water rotational process in a mixture and instead relax with a longer relaxation time, $\tau_c$, can be estimated from the dielectric strength, $\Delta \epsilon_{wp}$, of the corresponding relaxation process. The hydration number, which denotes the average number of water molecules in the hydration layer of a glycerol molecule, is given by 

$$N_{hyd}(c) = \frac{\epsilon_w - \Delta \epsilon_{wp}}{\Delta \epsilon_{wp}}$$

(5)

where $c$ is the glycerol molar concentration, $\epsilon_w$ is the molar concentration of water in the mixture, and $\epsilon_{wp} = 55.35$ M is the molarity and $\Delta \epsilon_{wp} = 73.25$ is the dielectric strength of pure water at 25°C.

The dielectric strength of bound water molecules varies nonlinearly with glycerol concentration, $x_{glyc}$. At low glycerol concentrations ($0 < x_{glyc} < 10$ mol %), the dielectric strength of bound water increases linearly with glycerol concentration. When the glycerol concentration is higher than about 10 mol %, the dielectric strength shows a saturation behavior. In the low-concentration regime, our analysis shows that by average, 5.58 water molecules are present in the hydration shell around a glycerol molecule, which agrees well with the average number of water molecules in the primary hydration layer of glycerol.
5.57, directly computed in MD simulations for low glycerol concentration solutions. This number is expected as each glycerol molecule can form at most six hydrogen bonds with surrounding water molecules through its three OH groups. The value also agrees with other reports in the literature.\textsuperscript{10,47} In a dilute mixture, glycerol molecules are uniformly dispersed in the mixture. The average number of water molecules bound to a glycerol molecule is roughly constant as long as the hydration shells of different glycerol molecules do not overlap. When the glycerol concentration is increased beyond a certain value, the hydration layers start to overlap and glycerol molecules aggregate to form clusters and networks, resulting in a decrease in the hydration number. Because the dielectric response of the bound water shows a saturation behavior at a specific glycerol concentration of \( \sim 7.5\) mol %, this concentration is roughly the threshold value signaling the overlapping of hydration shells. Observations of a similar transition of the hydration structure have been reported for aqueous solutions of bovine serum albumin,\textsuperscript{25} lysozyme proteins,\textsuperscript{27,43} and micelles.\textsuperscript{26}

### 3.3. Confined Water in the Glycerol Network

Water and glycerol molecules are well associated in mixtures with high glycerol concentrations. Dielectric spectra of high glycerol concentration solutions suggested that water cooperative domains do not exist and water molecules are dispersed well in the mixtures.\textsuperscript{17,22} The long relaxation time extracted from our dielectric response measurements (Figure 2), \( \tau_c = 85 \pm 9\) ps, corresponding to a characteristic frequency of 1.87 \( \pm 0.22\) GHz, is for water confined in a glycerol network and strongly bound to more than one glycerol molecule. In other words, these water molecules are in overlapped hydration shells. Because these water molecules have strong hydrogen bonds with more than one glycerol molecule, they relax much more slowly. Furthermore, this relaxation time is almost constant as the glycerol concentration is varied (Figure 5, inset). In a mixture with a low glycerol concentration, the contribution of the confined water to the overall dielectric response of the mixture is negligible, indicating that the amount of confined water is insignificant. When the glycerol concentration is higher than about \(7.5\) mol %, hydration shells of different glycerol molecules start to overlap and the confined water emerges. This value coincides with that observed for the saturation value of the dielectric strength of bound water in Figure 4b and the increase of the effective dipole moment of glycerol in solutions (Figure 3, lower inset). After that point, the amount of confined water grows almost linearly with the glycerol concentration, as indicated by the dielectric strength, \( \Delta \varepsilon_c \), of the corresponding relaxation process (Figure 5). In a glycerol-rich mixture, this category of water dominates and the dielectric response of the mixture is mainly controlled by the glycerol network and the water confined in it. This physical picture motivates us to perform MD simulations of glycerol–water mixtures and extract relaxation times of water as discussed below.

### 3.4. MD Simulations

To further probe the dynamics of water and glycerol at the molecular scales, we conducted all-atom MD simulation of their mixtures. All simulations were performed with a large-scale atomic/molecular massively parallel simulator (LAMMPS).\textsuperscript{48} Totally, seven systems with glycerol molar fraction ranging from 0 (pure water) to 1 (pure glycerol) were simulated, and their parameters are listed in Table 2. The structure of a glycerol molecule was constructed using Automated Topology Builder.\textsuperscript{49} The GROMOS force field\textsuperscript{50} was adopted for glycerol, and the SPC/E model\textsuperscript{14,52} was used for water. The cut-off of the 12-6 Lennard-Jones potential was set as 14 \( \AA \) for glycerol and 7.9 \( \AA \) for water. Coulomb interactions were fully accounted for with the long-range part computed with the particle–particle particle-mesh method. The geometric mixing rule was adopted for glycerol–water interactions. The equations of motions were integrated with a velocity Verlet algorithm with a time step of 1 fs. Each system was equilibrated in an NPT ensemble at 1 atmospheric pressure and 300 K for 10 ns. After the density of each system became stable, an NVT ensemble was used for production runs with the system temperature fixed at 300 K using a Nose–Hoover thermostat. For the pure glycerol system, the production run was 10 ns, and a snapshot of the system was dumped every ps. For all other systems, the production run was 500 ps, and a snapshot was dumped every 0.05 ps.

We calculated the relaxation autocorrelation function, \( C(t) \), of a dipole using eq 6 in which \( \mu(t) \) is the unit electric dipole associated with a water or a glycerol molecule at time \( t \), and the summation is over all \( N \) molecules (either water or glycerol) in a system. An ensemble average was computed with various states of the system being taken as the initial state at \( t = 0 \) to get the autocorrelation function

\[
C(t) = \left\langle \frac{1}{N} \sum_{i=1}^{N} \mu_i(t)\mu_i(0) \right\rangle
\]

An analysis indicates that all relaxation autocorrelation functions can be fit to a stretched exponential form, \( e^{-(t/\tau)^{\beta}} \), with \( \tau \) being average relaxation time and \( \beta \) being an exponent. When \( \beta = 1 \), the relaxation is exponential. The relaxation autocorrelation functions as well as their fits for pure water, five glycerol–water mixtures are shown in Figure 6a, and the results for pure glycerol is presented in the inset of Figure 6a. All results on \( \tau \) and \( \beta \) from such a fit are included in Table 2. For pure water, \( \beta \) is close to 1 and the relaxation time is about \( 4.9\) ps, which is very close to the value reported in the literature.\textsuperscript{53} For pure glycerol, \( \beta \) is about 0.69 and the relaxation time is about 986 ps. These relaxation time values

### Table 2. Glycerol–Water Mixtures Studied in MD Simulations

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<th>mol %</th>
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<th>( N_{\text{glyc}} )</th>
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\( \text{The first four columns show the molar fraction of glycerol, the number of water molecules, the number of glycerol molecules, and the size of the cubic simulation box for each system.} \) \( N_{\text{hyd}} \) is the number of bound water molecules per glycerol molecule, which are defined as the water molecules in the primary hydration shell of a glycerol molecule. \( f_s \) is the fraction of water in each system as the bound water and mathematically, \( f_s = N_{\text{hyd}} \times N_{\text{glyc}}/N_{\text{water}} \). The relaxation autocorrelation functions, \( C(t) \), are fitted to a stretched exponential form \( e^{-(t/\tau)^{\beta}} \) with \( \beta \) being an exponent and \( \tau \) being average relaxation time. Here, we show the average relaxation time of water in each solution except in pure glycerol, where \( \tau \) indicates the relaxation time of glycerol.

\[ \text{C}(t) = \left\langle \frac{1}{N} \sum_{i=1}^{N} \mu_i(t)\mu_i(0) \right\rangle \]

\[ \frac{\tau}{\tau_{\text{water}}} = \left( \frac{\rho_{\text{water}}}{\rho_{\text{glycerol}}} \right)^{1/\beta} \]

\[ \frac{\rho_{\text{glycerol}}}{\rho_{\text{water}}} = \left( \frac{\tau_{\text{water}}}{\tau} \right)^{1/\beta} \]
and 35 mol %. The relaxation autocorrelation functions, analysis of mixtures with a high glycerol concentration of 19.69 mol %, shows a deviation from a linear dependence (dash line) on the linear scale for this mixture (inset of Figure 6). On short timescales, the relaxation is dominated by a relaxation time of about 28 ps (50 ps). When the glycerol concentration is increased, the average relaxation time has increased from 0 to 10 mol %, the average relaxation time has ∼8 ps for bulk water and 1100 ps for glycerol. However, on longer timescales, the relaxation is almost exponential as indicated by the linear decay of $C(t)$ in the region of $30 \text{ ps} < t < 100 \text{ ps}$. The corresponding relaxation time is about 28 ps, which is only slightly lower than the experimental value of 35 ps identified from the dielectric spectroscopy data. The difference can be partially attributed to the SPC/E model adopted here for water. For bulk water, a similar trend can be noted that the relaxation time from the SPC/E model is 4.9 ps, which is also lower than the experimental value of ∼8 ps.

To better understand the structure of hydration shells around glycerol molecules, we analyze the distribution of water molecules when the glycerol concentration is varied. For this study, we consider the position of the oxygen atom in each water molecule as the center point of that molecule. At a given time, we compute the shortest distance, $r$, between a water molecule and an oxygen or carbon atom on any glycerol molecule in a mixture. We then count the number of water molecules in glycerol hydration shells from $r - \delta r$ to $r + \delta r$ and normalize this number by the number of glycerol molecules, $N_{\text{glyc}}$, in the system. The result of this analysis is denoted as the normalized hydration number, $N(r)$, as a function of distance, $r$, at various glycerol concentrations (Figure 7). Two main zones including an excluded zone and a primary hydration layer around a glycerol molecule are clearly shown in these functions. The excluded zone extends from the surface of a glycerol molecule, where $r = 0$, to $r \approx 2.3 \text{ Å}$ and the primary hydration layer extends to $r \approx 3.15 \text{ Å}$, indicating a thickness of about 0.85 Å for the hydration layer. All water molecules in the primary hydration shell are referred to as the bound water and denoted as the hydration number, $N_{\text{hyd}}$. The ratio $f_b = N_{\text{hyd}} / N_{\text{glyc}}$ is $N_{\text{water}} / N_{\text{water}}$ with $N_{\text{water}}$ being the total number of water molecules in a solution, thus, indicates the fraction of water located in primary hydration shells of glycerol molecules. The results on $N_{\text{hyd}}$ and $f_b$ for all systems modeled here are included in Table 2. Specifically, for the mixture containing 4.15 mol % glycerol, $N_{\text{hyd}}$ is around 5.57, which matches well with the experimental value of 5.58 water molecules in the hydration layer of a glycerol molecule in the low-concentration regime. When the glycerol concentration is increased, $N_{\text{hyd}}$...
gradually decreases and at the glycerol concentration of 50 mol %, the fraction of water being bound to glycerol, $f_w$, gradually increases toward 1 (Figure 6b, inset), meaning that, by average, there is one water molecule per one glycerol molecule in this concentrated mixture. Furthermore, Figure 7 indicates that when the glycerol concentration is increased beyond 19 mol %, the second peak in the normalized hydration number curve disappears. This behavior is consistent with the physical picture that at high glycerol concentrations, the hydration shells of adjacent glycerol molecules start to overlap.

When the hydration shells start to overlap, a much longer relaxation time emerges in our simulations. As shown in Figure 6b, the average relaxation time dramatically increases and shows a clear deviation from a linear dependence on the glycerol concentration when it is beyond about 10 mol %. In highly concentrated mixtures, water molecules are strongly bound to glycerol molecules; thus, the average relaxation time increase significantly. Specifically, when the glycerol concentration is increased to 50 mol % ($f_w = 0.97$), 97% of water molecules are confined in primary hydration shells (2.3 Å < r < 3.15 Å) of glycerol molecules. The hydration shells of glycerol molecules strongly overlap, and all water molecules are confined in a glycerol network. The average relaxation time of water in the mixture with 50 mol % of glycerol is about 58 ps. A similar relaxation time, 50 ps, can be extracted from the relaxation autocorrelation function for the mixture containing 35 mol % of glycerol (Figure 6b on the log−linear scale). Therefore, the relaxation time of confined water is somewhere between 50 and 60 ps from MD simulations. This value is lower than the value of 85 ps identified experimentally through dielectric spectroscopy for water confined in a glycerol network. Again, the difference is largely due to the fact that the SPC/E model generally underestimates the relaxation times of water.

4. CONCLUSIONS

We have performed the dielectric spectroscopy of glycerol−water mixtures in a wide frequency range from the megahertz- to terahertz region to systematically inspect the transition from pure water toward pure glycerol. An analysis of the dielectric response of glycerol−water mixtures has revealed four distinct relaxation processes including the rotational motion of glycerol molecules with a reorientational time of ~910 ps, water confined in a glycerol network with a relaxation time of ~85 ps, water bound in the hydration layer of a glycerol molecule with a relaxation time of ~35 ps, and bulk water with a relaxation time of ~8 ps. A critical glycerol concentration of ~7.5 mol % has been identified. Below this threshold, the dielectric response of the mixture is controlled by bulk water, bound water, and glycerol. Beyond the critical concentration, confined water emerges and contributes to the dielectric response of the mixture as well. In the regime of low glycerol concentrations, by average, the hydration shell of a glycerol molecule consists of about 5.58 water molecules. In mixtures with higher glycerol concentrations, the hydration shells start to merge and overlap, and the dielectric response from bound water shows a saturation behavior, whereas the dielectric response from confined water increases with an increasing glycerol concentration. The physical picture revealed from the dielectric spectroscopy is further confirmed with MD simulations. The results provide an in-depth understanding of the dynamics of water and glycerol in their mixtures and insights into the reactivity of glycerol as a common cosolvent of water.

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Notes

The authors declare no competing financial interest.

Acknowledgments

This material is based upon work supported by the Air Force Office of Scientific Research under award number FA9550-18-1-0263 and National Science Foundation (CHE-1665157). We acknowledge Advanced Research Computing at Virginia Tech for providing computational resources and technical support that have contributed to the results reported within this paper.

References


